

The Cahn-Hilliard-Oono equation

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The Cahn–Hilliard equation : recent advances and applications

$$\frac{\partial u}{\partial t} + \alpha \kappa \Delta^2 u - \kappa \Delta f(u) + \beta u = 0, \quad \beta > 0$$

Models long-ranged interactions

Based on the free energy :

$$\Psi_{\Omega} = \int_{\Omega} \left(\frac{\alpha}{2} |\nabla u|^2 + F(u) + \int_{\Omega} u(y) k(x, y) u(x) dy \right) dx$$

$$k(x, y) = \frac{\beta}{4\pi|x - y|}$$

Evolution equation :

$$\begin{aligned} \frac{\partial u}{\partial t} &= \kappa \Delta \partial_u \Psi_{\Omega} \\ \Delta(-k(x, y)) &= d_i(x - y) \end{aligned}$$

$$\frac{\partial u}{\partial t} + \beta u + \Delta^2 u - \Delta f(u) = 0, \beta > 0$$

$$\frac{\partial u}{\partial \nu} = \frac{\partial \Delta u}{\partial \nu} = 0 \text{ on } \Gamma$$

$$u|_{t=0} = u_0$$

No mass conservation :

$$\frac{d\langle u \rangle}{dt} + \beta \langle u \rangle = 0$$

$$\langle \cdot \rangle = \frac{1}{\text{Vol}(\Omega)} \int_{\Omega} \cdot dx \text{ or } \frac{1}{\text{Vol}(\Omega)} \langle \cdot, 1 \rangle_{H^{-1}(\Omega), H^1(\Omega)}$$

Thus :

$$\langle u(t) \rangle = e^{-\beta t} \langle u_0 \rangle, t \geq 0$$

→ Conservation of mass only when $\langle u_0 \rangle = 0$

Note that

$$|\langle u(t) \rangle| \leq |\langle u_0 \rangle|, \quad t \geq 0$$

$$\lim_{t \rightarrow +\infty} \langle u(t) \rangle = 0 \quad (\beta > 0 \text{ fixed})$$

$$\lim_{\beta \rightarrow 0^+} \langle u(t) \rangle = \langle u_0 \rangle \quad (t \geq 0 \text{ fixed})$$

Simplest equation for which we do not have the conservation of mass

Already brings several difficulties

Cubic nonlinear term :

$$f(s) = s^3 - s \text{ (can be generalized)}$$

Assume that $|\langle u_0 \rangle| \leq M, M \geq 0$

A priori :

$$|\langle u(t) \rangle| \leq M, \forall t \geq 0$$

A priori estimates :

Equivalent equation :

$$\frac{\partial \bar{u}}{\partial t} + \beta \bar{u} + \Delta^2 u - \Delta f(u) = 0$$

$$\bar{u} = u - \langle u \rangle$$

Weak formulation :

$$(-\Delta)^{-1} \frac{\partial \bar{u}}{\partial t} + \beta (-\Delta)^{-1} \bar{u} - \Delta u + \overline{f(u)} = 0$$

Multiply by \bar{u} :

$$\frac{1}{2} \frac{d}{dt} \|\bar{u}\|_{-1}^2 + \beta \|\bar{u}\|_{-1}^2 + \|\nabla u\|^2 + ((f(u), \bar{u})) = 0$$

Hölder's inequality :

$$((f(u), u)) \geq \frac{3}{4} \|u\|_{L^4(\Omega)}^4 - c$$

$$|\langle u \rangle \int_{\Omega} f(u) dx| \leq \frac{1}{4} \|u\|_{L^4(\Omega)}^4 + c_M$$

Thus :

$$\frac{d}{dt} \|\bar{u}\|_{-1}^2 + \|\nabla u\|^2 + \|u\|_{L^4(\Omega)}^4 \leq c_M$$

Finally :

$$\|\bar{u}\|_{-1} \leq c \|\nabla u\|,$$

Thus :

$$\frac{d}{dt} \|\bar{u}\|_{-1}^2 + c \|\bar{u}\|_{-1}^2 + \frac{1}{2} \|\nabla u\|^2 + \|u\|_{L^4(\Omega)}^4 \leq c'_M, \quad c > 0$$

All constants are independent of β

Note that :

$$\frac{d\langle u \rangle^2}{dt} = 2\langle u \rangle \frac{d\langle u \rangle}{dt} = -2\beta \langle u \rangle^2 \leq 0$$

Thus :

$$\frac{d}{dt} (\|\bar{u}\|_{-1}^2 + \langle u \rangle^2) + c (\|\bar{u}\|_{-1}^2 + \langle u \rangle^2) + \frac{1}{2} \|\nabla u\|^2 + \|u\|_{L^4(\Omega)}^4 \leq c'_M, \quad c > 0$$

→ Dissipative estimate on $\|\bar{u}\|_{-1}^2 + \langle u \rangle^2$:

$$\|\bar{u}(t)\|_{-1}^2 + \langle u(t) \rangle^2 \leq e^{-ct}(\|\bar{u}_0\|_{-1}^2 + \langle u_0 \rangle^2) + c', \quad c > 0$$

Bounded absorbing set in $H^{-1}(\Omega)$: $\forall R > 0$, $\|u_0\|_{H^{-1}(\Omega)} \leq R$,

$\exists t_0 = t_0(R) \geq 0$ such that $t \geq t_0$ implies

$$\|u(t)\|_{H^{-1}(\Omega)} \leq c$$

Furthermore :

$$\int_t^{t+r} \|\nabla u\|^2 ds \leq c_r, \quad t \geq t_0, \quad r > 0$$

$$\int_t^{t+r} \|u\|_{L^4(\Omega)}^4 ds \leq c_r, \quad t \geq t_0, \quad r > 0$$

Multiply by $-\Delta u$ ($f' \geq -1$) :

$$\frac{d}{dt} \|\bar{u}\|^2 + \|\Delta u\|^2 \leq 2\|\nabla u\|^2$$

Note that

$$\|u\|^2 \leq 2(\|\bar{u}\|^2 + \langle u \rangle^2) \leq c(\|\nabla u\|^2 + M^2),$$

Thus :

$$\begin{aligned} \|u(t)\| &\leq c, \quad t \geq t_1 (\geq t_0) \\ \int_t^{t+r} \|\Delta u\|^2 ds &\leq c_r, \quad t \geq t_1 \end{aligned}$$

→ Existence of a bounded absorbing set in $L^2(\Omega)$

Multiply by $\frac{\partial \bar{u}}{\partial t}$:

$$\frac{d}{dt}(\beta \|\bar{u}\|_{-1}^2 + \|\nabla u\|^2 + 2 \int_{\Omega} F(u) dx) + 2 \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 + 2 \text{Vol}(\Omega) \langle f(u) \rangle \left\langle \frac{\partial u}{\partial t} \right\rangle = 0$$

$$F(s) = \int_0^s f(\xi) d\xi$$

Young's inequality :

$$|2 \text{Vol}(\Omega) \langle f(u) \rangle \left\langle \frac{\partial u}{\partial t} \right\rangle| \leq c\beta M \left(\int_{\Omega} u^4 dx + 1 \right)$$

Thus :

$$\frac{d}{dt}(\beta \|\bar{u}\|_{-1}^2 + \|\nabla u\|^2 + 2 \int_{\Omega} F(u) dx) + \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 \leq c\beta M \left(\int_{\Omega} u^4 dx + 1 \right)$$

Uniform Gronwall's lemma :

$$\|u(t)\|_{H^1(\Omega)} \leq c, \quad t \geq t_2 (\geq t_1)$$

→ Existence of a bounded absorbing set in $H^1(\Omega)$

Upper bounds : depends on β (depends continuously on β)

We can obtain more regularity

The dissipative semigroup :

Existence : follows from the a priori estimates

Uniqueness :

Let u_1, u_2 be two solutions with initial data $u_{1,0}, u_{2,0}$

Set $u = u_1 - u_2, u_0 = u_{1,0} - u_{2,0}$:

$$\frac{\partial u}{\partial t} + \beta u + \Delta^2 u - \Delta(f(u_1) - f(u_2)) = 0$$

$$\frac{\partial u}{\partial \nu} = \frac{\partial \Delta u}{\partial \nu} = 0 \text{ on } \Gamma$$

$$u|_{t=0} = u_0$$

Equivalent formulation :

$$\frac{\partial \bar{u}}{\partial t} + \beta \bar{u} + \Delta^2 u - \Delta(f(u_1) - f(u_2)) = 0$$

Multiply by $(-\Delta)^{-1}\bar{u}$:

$$\frac{1}{2} \frac{d}{dt} \|\bar{u}\|_{-1}^2 + \beta \|\bar{u}\|_{-1}^2 + \|\nabla u\|^2 + ((f(u_1) - f(u_2), \bar{u})) = 0$$

Note that

$$\begin{aligned} ((f(u_1) - f(u_2), \bar{u})) &\geq -\|u\|^2 - \langle u \rangle \int_{\Omega} (f(u_1) - f(u_2)) dx \\ \|u\|^2 &\leq 2(\|\bar{u}\|^2 + \langle u \rangle^2) \leq c(\|\bar{u}\|_{-1} \|\nabla u\| + \langle u \rangle^2) \\ &\leq \gamma \|\nabla u\|^2 + c(\|\bar{u}\|_{-1}^2 + \langle u \rangle^2), \quad \forall \gamma > 0 \end{aligned}$$

Thus :

$$\begin{aligned} |\langle u \rangle \int_{\Omega} (f(u_1) - f(u_2)) dx| &\leq c |\langle u \rangle| \int_{\Omega} (|u_1|^2 + |u_2|^2 + 1) |u| dx \\ &\leq c(\|u\|^2 + (\|u_1\|_{L^4(\Omega)}^4 + \|u_2\|_{L^4(\Omega)}^4 + 1) \langle u \rangle^2) \\ &\leq \frac{1}{4} \|\nabla u\|^2 + c(\|u_1\|_{L^4(\Omega)}^4 + \|u_2\|_{L^4(\Omega)}^4 + 1)(\|\bar{u}\|_{-1}^2 + \langle u \rangle^2) \end{aligned}$$

Thus :

$$\frac{d}{dt} \|\bar{u}\|_{-1}^2 + \|\nabla u\|^2 \leq c(\|u_1\|_{L^4(\Omega)}^4 + \|u_2\|_{L^4(\Omega)}^4 + 1)(\|\bar{u}\|_{-1}^2 + \langle u \rangle^2)$$

Note that $\frac{d\langle u \rangle^2}{dt} \leq 0$:

$$\frac{d}{dt} (\|\bar{u}\|_{-1}^2 + \langle u \rangle^2) + \|\nabla u\|^2 \leq c(\|u_1\|_{L^4(\Omega)}^4 + \|u_2\|_{L^4(\Omega)}^4 + 1)(\|\bar{u}\|_{-1}^2 + \langle u \rangle^2)$$

Gronwall's lemma :

$$\|u(t)\|_{H^{-1}(\Omega)} \leq ce^{c't} \|u_0\|_{H^{-1}(\Omega)}$$

The constants are independent of β

→ Continuous dependence and uniqueness

We can define the continuous semigroups :

$$S_\beta(t) : L^2(\Omega) \rightarrow L^2(\Omega), u_0 \mapsto u(t), t \geq 0, \beta > 0$$

$$S_\beta(t) : L^2(\Omega) \rightarrow H^1(\Omega), t > 0$$

Set :

$$\Phi_M = \{v \in L^2(\Omega), |\langle v \rangle| \leq M\}, M \geq 0$$

$$S_\beta(t) : \Phi_M \rightarrow \Phi_M, t \geq 0$$

Theorem : The semigroup $S_\beta(t)$ possesses the finite-dimensional (for the $H^{-1}(\Omega)$ -topology) global attractor \mathcal{A}_β^M on the phase space Φ_M which is compact in $L^2(\Omega)$ and bounded in $H^1(\Omega)$.

Finite fractal dimensionality :

Multiply

$$\frac{\partial u}{\partial t} + \beta u + \Delta^2 u - \Delta(f(u_1) - f(u_2)) = 0$$

by tu :

$$\frac{d}{dt}(t\|u\|^2) \leq \|u\|^2 + ct\|\nabla u\|^2$$

Thus :

$$\|u(t)\|^2 \leq c \frac{1+t}{t} \int_0^t \|u\|_{H^1(\Omega)}^2 ds, \quad t > 0$$

Furthermore :

$$\int_0^t \|\nabla u\|^2 ds \leq c e^{c't} \|u_0\|_{H^{-1}(\Omega)}^2$$

$$\|u\|^2 \leq c(\|\nabla u\|^2 + \langle u \rangle^2)$$

$$\langle u \rangle^2 \leq \langle u_0 \rangle^2$$

Thus :

$$\|u(t)\|^2 \leq c \frac{1+t}{t} e^{c't} \|u_0\|_{H^{-1}(\Omega)}^2$$

All constants are independent of β

Remarks :

a) The bounds on, as well as the upper bound on the fractal dimension of, \mathcal{A}_β^M are independent of β , $\beta \leq \beta_0$

b) We can extend the semigroups to

$$\tilde{\Phi}_M = \{v \in H^{-1}(\Omega), |\langle v \rangle| \leq M\}, M \geq 0$$

Note that

$$S_\beta(t) : \tilde{\Phi}_M \rightarrow \Phi_M, t > 0$$

Definition : A compact set $\mathcal{M} \subset E$ is an exponential attractor for the semigroup $S(t)$ acting on the Banach space E if

- (i) It has finite fractal dimension, $\dim_F \mathcal{M} < +\infty$.
- (ii) It is positively invariant, $S(t)\mathcal{M} \subset \mathcal{M}$, $\forall t \geq 0$.
- (iii) It attracts exponentially fast the bounded subsets of E in the following sense :

$$\forall B \subset E \text{ bounded, } \text{dist}_E(S(t)B, \mathcal{M}) \leq Q(\|B\|_E)e^{-\alpha t}, \quad t \geq 0,$$

where the positive constant α and the function Q are independent of B .

Construction of robust exponential attractors

Let E and E_1 be two Banach spaces such that the embedding $E_1 \subset E$ is compact, X be a bounded subset of E and $L_\epsilon : X \rightarrow X$, $\epsilon \in [0, \epsilon_0]$, $\epsilon_0 > 0$, be a family of operators such that

1. For every $x_1, x_2 \in X$ and every $\epsilon \in [0, \epsilon_0]$,

$$\|L_\epsilon x_1 - L_\epsilon x_2\|_{E_1} \leq c \|x_1 - x_2\|_E,$$

where the constant c is independent of ϵ .

2. For every $\epsilon \in [0, \epsilon_0]$, every $i \in \mathbb{N}$ and every $x \in X$,

$$\|L_\epsilon^i x - L_0^i x\|_E \leq c^i \epsilon,$$

where the constant c is independent of ϵ .

Then, there exists a family $\mathcal{M}_\epsilon \subset X$, $\epsilon \in [0, \epsilon_0]$, such that \mathcal{M}_ϵ is an exponential attractor for the discrete dynamical system generated by L_ϵ , i.e.,

(i) The set \mathcal{M}_ϵ is compact and has finite fractal dimension in E ,

$$\dim_F \mathcal{M}_\epsilon \leq c.$$

(ii) The set \mathcal{M}_ϵ is positively invariant,

$$L_\epsilon \mathcal{M}_\epsilon \subset \mathcal{M}_\epsilon.$$

(iii) The set \mathcal{M}_ϵ attracts X exponentially fast,

$$\text{dist}_E(L^i X, \mathcal{M}_\epsilon) \leq ce^{-c'i}, \quad i \in \mathbb{N}, \quad c' > 0.$$

(iv) The family \mathcal{M}_ϵ is Hölder continuous at $\epsilon = 0$,

$$\text{dist}_{\text{sym}}(\mathcal{M}_\epsilon, \mathcal{M}_0) \leq c\epsilon^{c'}, \quad c' \in (0, 1),$$

where dist_{sym} denotes the Hausdorff symmetric distance between sets defined as

$$\text{dist}_{\text{sym}}(A, B) = \max(\text{dist}_E(A, B), \text{dist}_E(B, A)).$$

Finally, all constants are independent of ϵ and can be computed explicitly.

Theorem : For every $\beta \in [0, \beta_0]$, $\beta_0 > 0$ given, the semigroup $S_\beta(t)$ acting on $\tilde{\Phi}_M$ possesses an exponential attractor \mathcal{M}_β^M on $\tilde{\Phi}_M$ such that

1. The set \mathcal{M}_β^M has finite fractal dimension in $H^{-1}(\Omega)$,

$$\dim_{\text{F}} \mathcal{M}_\beta^M \leq c.$$

2. The set \mathcal{M}_β^M is positively invariant by $S_\beta(t)$,

$$S_\beta(t) \mathcal{M}_\beta^M \subset \mathcal{M}_\beta^M, \quad t \geq 0.$$

3. The set \mathcal{M}_β^M attracts all bounded subsets of $\tilde{\Phi}_M$ exponentially fast, i.e., for every bounded subset B of $\tilde{\Phi}_M$, there exists a constant $c = c(B)$ such that

$$\text{dist}_{H^{-1}(\Omega)}(S_\beta(t)B, \mathcal{M}_\beta^M) \leq ce^{-c't}, \quad t \geq 0, \quad c' > 0.$$

4. The family of sets \mathcal{M}_β^M is Hölder continuous at 0,

$$\text{dist}_{\text{sym}}(\mathcal{M}_\beta^M, \mathcal{M}_0^M) \leq c\beta^{c'}, \quad c' \in (0, 1).$$

Furthermore, all constants are independent of β and can be computed explicitly.

Existence of a uniform (with respect to β) bounded absorbing set $B_0 \subset \Phi_M \cap H^1(\Omega)$, i.e., $\forall B \subset \Phi_M$ bounded, $\exists t_0 = t_0(B) > 0$ independent of $\beta \in [0, \beta_0]$ such that

$$S_\beta(t)B \subset B_0, \quad t \geq t_0, \quad \beta \in [0, \beta_0]$$

→ It suffices to construct the exponential attractor \mathcal{M}_β^M on B_0

There exists $t_1 > 0$ independent of $\beta \in [0, \beta_0]$ such that

$$S_\beta(t)B_0 \subset B_0, \quad t \geq t_1, \quad \beta \in [0, \beta_0]$$

Let u^β, u^0 be two solutions to the problem for $\beta > 0, \beta = 0$ with the same initial datum u_0

Set $u = u^\beta - u^0$:

$$\frac{\partial u}{\partial t} + \beta u + \Delta^2 u - \Delta(f(u^\beta) - f(u^0)) = -\beta u^0$$

$$\frac{\partial u}{\partial \nu} = \frac{\partial \Delta u}{\partial \nu} = 0 \text{ on } \Gamma$$

$$u|_{t=0} = 0$$

Recall that

$$\frac{d\langle u^\beta \rangle}{dt} + \beta \langle u^\beta \rangle = 0 \quad (\beta > 0)$$

$$\frac{d\langle u^0 \rangle}{dt} = 0$$

Thus :

$$\frac{\partial \bar{u}}{\partial t} + \beta \bar{u} + \Delta^2 u - \Delta(f(u^\beta) - f(u^0)) = -\beta \bar{u}^0$$

Multiply by $(-\Delta)^{-1} \bar{u}$:

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|\bar{u}\|_{-1}^2 + \beta \|\bar{u}\|_{-1}^2 + \|\nabla u\|^2 + ((f(u^\beta) - f(u^0), \bar{u})) \\ = -\beta (((-\Delta)^{-\frac{1}{2}} \bar{u}^0, (-\Delta)^{-\frac{1}{2}} \bar{u})) \end{aligned}$$

Recall that

$$\|(-\Delta)^{-\frac{1}{2}}\bar{u}\| \leq c\|\nabla u\|$$

Thus :

$$\frac{d}{dt}\|\bar{u}\|_{-1}^2 + \|\nabla u\|^2 \leq c\beta^2\|\bar{u}^0\|_{-1}^2 + 2\|u\|^2 + 2\langle u \rangle \int_{\Omega} (f(u^\beta) - f(u^0)) dx$$

Recall that

$$\begin{aligned}\|u\|^2 &\leq 2(\|\bar{u}\|^2 + \langle u \rangle^2) \leq c(\|\bar{u}\|_{-1}\|\nabla u\| + \langle u \rangle^2) \\ &\leq \gamma\|\nabla u\|^2 + c(\|\bar{u}\|_{-1}^2 + \langle u \rangle^2), \quad \forall \gamma > 0\end{aligned}$$

Thus :

$$\begin{aligned}|2\langle u \rangle \int_{\Omega} (f(u^\beta) - f(u^0)) dx| &\leq c|\langle u \rangle| \int_{\Omega} (|u^\beta|^2 + |u^0|^2 + 1)|u| dx \\ &\leq c(\|u\|^2 + (\|u^\beta\|_{L^4(\Omega)}^4 + \|u^0\|_{L^4(\Omega)}^4 + 1)\langle u \rangle^2)\end{aligned}$$

$$\leq \frac{1}{2} \|\nabla u\|^2 + c(\|u^\beta\|_{L^4(\Omega)}^4 + \|u^0\|_{L^4(\Omega)}^4 + 1)(\|\bar{u}\|_{-1}^2 + \langle u \rangle^2)$$

Thus :

$$\frac{d}{dt} \|\bar{u}\|_{-1}^2 \leq c\beta^2 \|\bar{u}^0\|_{-1}^2 + c'(\|u^\beta\|_{L^4(\Omega)}^4 + \|u^0\|_{L^4(\Omega)}^4 + 1)(\|\bar{u}\|_{-1}^2 + \langle u \rangle^2)$$

Finally :

$$\begin{aligned} \frac{d\langle u \rangle^2}{dt} &= 2\langle u \rangle(-\beta\langle u \rangle - \beta\langle u^0 \rangle) \leq -2\beta\langle u \rangle\langle u^0 \rangle \\ &\leq \beta^2\langle u^0 \rangle^2 + \langle u \rangle^2 \end{aligned}$$

Thus :

$$\begin{aligned} \frac{d}{dt}(\|\bar{u}\|_{-1}^2 + \langle u \rangle^2) &\leq c\beta^2(\|\bar{u}^0\|_{-1}^2 + \langle u^0 \rangle^2) \\ &+ c'(\|u^\beta\|_{L^4(\Omega)}^4 + \|u^0\|_{L^4(\Omega)}^4 + 1)(\|\bar{u}\|_{-1}^2 + \langle u \rangle^2) \end{aligned}$$

We have :

$$\int_0^t \|u^\beta\|_{L^4(\Omega)}^4 ds \leq ce^{c't}$$

$$\int_0^t \|u^0\|_{L^4(\Omega)}^4 ds \leq ce^{c't}$$

$$\|u^0\|_{H^{-1}(\Omega)} \leq c$$

Gronwall's lemma :

$$\|u(t)\|_{H^{-1}(\Omega)} \leq c\beta e^{c't}$$

The constants are independent of β

Set $L_\beta = S_\beta(t_1)$, take $E = H^{-1}(\Omega)$, $E_1 = L^2(\Omega)$

→ Existence of a robust family of exponential attractors $\mathcal{M}_\beta^{M,d}$ for the discrete dynamical systems generated by L_β

Set

$$\mathcal{M}_\beta^M = \cup_{t \in [0, t_1]} S_\beta(t) \mathcal{M}_\beta^{M, d}$$

We need to prove that the mapping $(t, x) \mapsto S_\beta(t)x$ is Hölder continuous on $[0, t_1] \times B_0$, uniformly with respect to $\beta \in [0, \beta_0]$

Hölder (Lipschitz) continuity with respect to x : continuous dependence estimate

Hölder continuity with respect to t :

$$\begin{aligned} \|S_\beta(t+s)u_0 - S_\beta(t)u_0\|_{H^{-1}(\Omega)} &\leq |s|^{\frac{1}{2}} \left(\int_t^{t+s} \left\| \frac{\partial u}{\partial t} \right\|_{H^{-1}(\Omega)}^2 d\tau \right)^{\frac{1}{2}} \\ &\leq c|s|^{\frac{1}{2}} \left(\int_t^{t+s} (\|u\|_{H^1(\Omega)}^2 + \|f(u)\|^2) d\tau \right)^{\frac{1}{2}} \end{aligned}$$

Logarithmic nonlinear terms :

Equations :

$$\frac{\partial u}{\partial t} + \beta u = \Delta \mu$$

$$\mu = -\Delta u + f(u)$$

$$\frac{\partial u}{\partial \nu} = \frac{\partial \mu}{\partial \nu} = 0 \text{ on } \Gamma$$

$$u|_{t=0} = u_0$$

$$f(s) = -\theta_c s + \frac{\theta}{2} \ln \frac{1+s}{1-s}, \quad s \in (-1, 1)$$

Recall that ($c_0 = \theta_c$)

$$f' \geq -c_0, \quad c_0 > 0$$

$$f(s) = f_1(s) - c_0 s, \quad f'_1 \geq 0$$

Set $F(s) = \int_0^s f(\xi) d\xi$

Approximated functions $f_N \in \mathcal{C}^1(\mathbb{R})$:

$$f_N(s) = f(-1 + \frac{1}{N}) + f'(-1 + \frac{1}{N})(s + 1 - \frac{1}{N}), \quad s < -1 + \frac{1}{N}$$

$$f_N(s) = f(s), \quad |s| \leq 1 - \frac{1}{N}$$

$$f_N(s) = f(1 - \frac{1}{N}) + f'(1 - \frac{1}{N})(s - 1 + \frac{1}{N}), \quad s > 1 - \frac{1}{N}$$

Recall that

$$f'_N \geq -c_0$$

For N large enough, $F_N(s) = \int_0^s f_N(\xi) d\xi$:

$$-c_1 \leq F_N(s) \leq c_2 f_N(s)s + c_3, \quad c_2 > 0, \quad c_1, c_2 \geq 0, \quad s \in \mathbb{R}$$

$$f_N(s)s \geq c_4 |f_N(s)| - c_5, \quad c_4 > 0, \quad c_5 \geq 0$$

$c_i, i = 1, \dots, 5$: independent of N

$$f_N(s)(s - m) \geq c_m(|f_N(s)| + F_N(s)) - c'_m,$$

$$c_m > 0, c'_m \geq 0, s \in \mathbb{R}, m \in (-1, 1)$$

c_m, c'_m : independent of N and depend continuously on m .

Approximated problems :

$$\frac{\partial u_N}{\partial t} + \beta u_N = \Delta \mu_N$$

$$\mu_N = -\Delta u_N + f_N(u_N)$$

$$\frac{\partial u_N}{\partial \nu} = \frac{\partial \mu_N}{\partial \nu} = 0 \text{ on } \Gamma$$

$$u_N|_{t=0} = u_0$$

A priori estimates :

Crucial step : a priori estimate independent of N on $f_N(u_N)$ in $L^2((0, T) \times \Omega)$,
 $T > 0$

Assume that $-1 < u_0(x) < 1$, a.e. x , and $|\langle u_0 \rangle| \leq 1 - \delta$, $\delta \in (0, 1)$
 δ : fixed constant

Equation for the spatial average :

$$\frac{d\langle u_N \rangle}{dt} + \beta \langle u_N \rangle = 0$$

Thus :

$$\langle u_N(t) \rangle = e^{-\beta t} \langle u_0 \rangle, \quad t \geq 0$$

$$|\langle u_N(t) \rangle| \leq |\langle u_0 \rangle|, \quad t \geq 0$$

$$|\langle u_N(t) \rangle| \leq 1 - \delta, \quad t \geq 0$$

Equivalent formulation :

$$\begin{aligned}\frac{\partial \bar{u}_N}{\partial t} + \beta \bar{u}_N &= \Delta \mu_N \\ \mu_N &= -\Delta \bar{u}_N + f_N(u_N)\end{aligned}$$

Multiply the first equation by μ_N :

$$\|\nabla \mu_N\|^2 + ((\frac{\partial \bar{u}_N}{\partial t}, \mu_N)) + \beta((\bar{u}_N, \mu_N)) = 0$$

Second equation :

$$\begin{aligned}((\frac{\partial \bar{u}_N}{\partial t}, \mu_N)) + \beta((\bar{u}_N, \mu_N)) &= \frac{1}{2} \frac{d}{dt} \|\bar{u}_N\|_V^2 + \frac{d}{dt} \int_{\Omega} F_N(u_N) dx \\ &\quad - ((\frac{d\langle u_N \rangle}{dt}, f_N(u_N))) + \beta \|\bar{u}_N\|_V^2 + \beta((\bar{u}_N, f_N(u_N)))\end{aligned}$$

Note that

$$-((\frac{d\langle u_N \rangle}{dt}, f_N(u_N))) = \beta((\langle u_N \rangle, f_N(u_N)))$$

Thus :

$$\begin{aligned} & \frac{d}{dt}(\|\bar{u}_N\|_V^2 + 2 \int_{\Omega} F_N(u_N) dx) \\ & + c(\|\bar{u}_N\|_V^2 + \|f_N(u_N)\|_{L^1(\Omega)} + \|\nabla \mu_N\|^2) \leq c', \quad c > 0 \end{aligned}$$

Multiply the second equation by \bar{u}_N :

$$\begin{aligned} \|\bar{u}_N\|_V^2 + c(\|f_N(u_N)\|_{L^1(\Omega)} + \int_{\Omega} F_N(u_N) dx) & \leq c' + ((\mu_N, \bar{u}_N)) = c' + ((\bar{\mu}_N, \bar{u}_N)) \\ & \leq c' + c'' \|\bar{u}_N\|_V \|\nabla \mu_N\| \end{aligned}$$

Thus :

$$\begin{aligned} \|\bar{u}_N\|_V^2 + c(\|f_N(u_N)\|_{L^1(\Omega)} + \int_{\Omega} F_N(u_N) dx) \\ \leq c' + c''(\|\nabla \mu_N\|^2 + \|\bar{u}_N\|_V^2), \quad c > 0 \end{aligned}$$

Combine the two estimates :

$$\frac{dE_N}{dt} + c(E_N + \|f_N(u_N)\|_{L^1(\Omega)} + \|\nabla \mu_N\|^2) \leq c', \quad c > 0$$

$$E_N = \|\bar{u}_N\|_V^2 + 2 \int_{\Omega} F_N(u_N) dx$$

Equivalent equation :

$$(-\Delta)^{-1} \frac{\partial \bar{u}_N}{\partial t} + \beta (-\Delta)^{-1} \bar{u}_N = -\bar{\mu}_N$$

Multiply by $\frac{\partial \bar{u}_N}{\partial t}$:

$$\frac{\beta}{2} \frac{d}{dt} \|\bar{u}_N\|_{-1}^2 + \left\| \frac{\partial \bar{u}_N}{\partial t} \right\|_{-1}^2 = -\left((\bar{\mu}_N, \frac{\partial \bar{u}_N}{\partial t}) \right)$$

Note that

$$\begin{aligned} |((\bar{\mu}_N, \frac{\partial \bar{u}_N}{\partial t}))| &\leq \|\frac{\partial \bar{u}_N}{\partial t}\|_{-1} \|\nabla \mu_N\| \\ &\leq \frac{1}{4} \|\frac{\partial \bar{u}_N}{\partial t}\|_{-1}^2 + c \|\nabla \mu_N\|^2 \end{aligned}$$

Thus :

$$\beta \frac{d}{dt} \|\bar{u}_N\|_{-1}^2 + \|\frac{\partial \bar{u}_N}{\partial t}\|_{-1}^2 \leq c \|\nabla \mu_N\|^2$$

Note that

$$\bar{\mu}_N = -\Delta \bar{u}_N + \overline{f_N(u_N)}$$

Multiply by $-\Delta \bar{u}_N$:

$$\begin{aligned} \|\Delta \bar{u}_N\|^2 &\leq c_0 \|\bar{u}_N\|_V^2 - ((\bar{\mu}_N, \Delta \bar{u}_N)) \\ &\leq c_0 \|\bar{u}_N\|_V^2 + \frac{1}{2} \|\Delta \bar{u}_N\|^2 + c \|\nabla \mu_N\|^2 \end{aligned}$$

Thus :

$$\|\Delta \bar{u}_N\|^2 \leq 2c_0 \|\bar{u}_N\|_V^2 + c \|\nabla \mu_N\|^2$$

Combine with the previous estimate :

$$\frac{dE_N}{dt} + c(E_N + \|\bar{u}_N\|_{H^2(\Omega)}^2 + \|f_N(u_N)\|_{L^1(\Omega)} + \|\nabla \mu_N\|^2) \leq c', \quad c > 0$$

Note that

$$\|\overline{f_N(u_N)}\| \leq c(\|\bar{u}_N\|_{H^2(\Omega)} + \|\nabla \mu_N\|)$$

Thus :

$$\begin{aligned} \frac{dE_N}{dt} + c(E_N + \|\bar{u}_N\|_{H^2(\Omega)}^2 + \|f_N(u_N)\|_{L^1(\Omega)} + \|\overline{f_N(u_N)}\|^2 + \|\nabla \mu_N\|^2) \\ \leq c', \quad c > 0 \end{aligned}$$

Finally :

$$\int_{\Omega} |f_N(u_N)| dx \leq c \left| \int_{\Omega} f_N(u_N) \bar{u}_N dx \right| + c$$

Thus :

$$|\langle f_N(u_N) \rangle| \leq c \|\bar{u}_N\| \|\overline{f_N(u_N)}\| + c'$$

Note that

$$\|f_N(u)\|^2 \leq c(\|\overline{f_N(u_N)}\|^2 + |\langle f_N(u_N) \rangle|^2)$$

Thus :

$$\begin{aligned} & \|f_N(u_N)\|_{L^2(\Omega \times (0,T))} \\ & \leq c(\|\overline{f_N(u_N)}\|_{L^2(\Omega \times (0,T))} + \|\bar{u}_N\|_{L^\infty(0,T;L^2(\Omega))} \|\overline{f_N(u_N)}\|_{L^2(\Omega \times (0,T))}) \\ & \quad + c', \quad T > 0 \end{aligned}$$

Poincaré's inequality :

$$\|\overline{f_N(u_N)}\|_{L^2(\Omega \times (0,T))}^2 \leq cE_N(0) + c'T$$

$$\|\bar{u}_N\|_{L^\infty(0,T;L^2(\Omega))}^2 \leq c\|\bar{u}_N\|_{L^\infty(0,T;V)}^2 \leq c'E_N(0) + c''T$$

Thus :

$$\|f_N(u_N)\|_{L^2(\Omega \times (0,T))} \leq c_{T,\delta}(\|u_0\|_{H^1(\Omega)}^2 + 1), \quad T > 0$$

Furthermore :

$$E_N(t) \leq e^{-ct}E_N(0) + c', \quad c > 0, \quad t \geq 0$$

$$\|u_N(t)\|_{H^1(\Omega)} \leq c_\delta e^{-c't}(\|u_0\|_{H^1(\Omega)} + 1) + c'', \quad c' > 0, \quad t \geq 0$$

$$\|\mu_N\|_{L^2(0,T;H^1(\Omega))} \leq c_{T,\delta}(\|u_0\|_{H^1(\Omega)}^2 + 1), \quad T > 0$$

The dissipative semigroup :

Theorem : We assume that u_0 is given such that $u_0 \in H^1(\Omega)$ and $-1 < u_0(x) < 1$, a.e. x , with $|\langle u_0 \rangle| < 1$. Then, the problem possesses a unique (weak) solution such that, $\forall T > 0$,

$$u \in \mathcal{C}([0, T]; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega)),$$

$$\frac{\partial u}{\partial t} \in L^2(0, T; H^1(\Omega)'),$$

$$\mu \in L^2(0, T; H^1(\Omega)).$$

Furthermore, $-1 < u(x, t) < 1$, a.e. (x, t) .

Set

$$\Phi_M = \{v \in H^1(\Omega) \cap L^\infty(\Omega), -1 < v(x) < 1, \text{ a.e. } x \in \Omega, \langle v \rangle = M\},$$

$$M \in (-1, 1)$$

We can define the continuous (for the $H^{-1}(\Omega)$ -norm) semigroup

$$S_\beta(t) : \Phi_M \rightarrow \Phi_M, u_0 \mapsto u(t), t \geq 0$$

$S_\beta(t)$ is dissipative in Φ_M

Regularity and separation from the pure states :

Proposition : The solution u satisfies

$$\frac{\partial u}{\partial t} \in L^\infty(r, +\infty; H^{-1}(\Omega)) \cap L^2(r, T; H^1(\Omega)),$$

$\forall r < T, r > 0$ and $T > 0$ given.

Equation :

$$(-\Delta)^{-1} \frac{\partial \bar{u}}{\partial t} + \beta(-\Delta)^{-1} \bar{u} = -\bar{\mu}$$

Differentiate with respect to time :

$$\begin{aligned} (-\Delta)^{-1} \frac{\partial}{\partial t} \frac{\partial \bar{u}}{\partial t} + \beta(-\Delta)^{-1} \frac{\partial \bar{u}}{\partial t} &= -\frac{\partial \bar{\mu}}{\partial t} \\ \frac{\partial \bar{\mu}}{\partial t} &= -\Delta \frac{\partial \bar{u}}{\partial t} + \overline{f'(u) \frac{\partial u}{\partial t}} \end{aligned}$$

Multiply by $\frac{\partial \bar{u}}{\partial t}$:

$$\frac{1}{2} \frac{d}{dt} \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 + \beta \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 + \left\| \frac{\partial \bar{u}}{\partial t} \right\|_V^2 + \left((f'(u) \frac{\partial u}{\partial t}, \frac{\partial \bar{u}}{\partial t}) \right) = 0$$

Furthermore :

$$\left((f'(u) \frac{\partial u}{\partial t}, \frac{\partial \bar{u}}{\partial t}) \right) \geq -c_0 \left\| \frac{\partial \bar{u}}{\partial t} \right\|^2 + \left((f'(u) \frac{d\langle u \rangle}{dt}, \frac{\partial \bar{u}}{\partial t}) \right)$$

Thus :

$$\frac{d}{dt} \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 + 2\beta \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 + \left\| \frac{\partial \bar{u}}{\partial t} \right\|_V^2 \leq c \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 - 2 \left((f'(u) \frac{d\langle u \rangle}{dt}, \frac{\partial \bar{u}}{\partial t}) \right)$$

Difficulty : handle the term

$$((f'(u) \frac{d\langle u \rangle}{dt}, \frac{\partial \bar{u}}{\partial t})) = -\beta \text{Vol}(\Omega) \langle u \rangle \langle f'(u) \frac{\partial \bar{u}}{\partial t} \rangle$$

To avoid it :

$$\begin{aligned} ((f'(u) \frac{\partial u}{\partial t}, \frac{\partial \bar{u}}{\partial t})) &= ((f'(u) \frac{\partial u}{\partial t}, \frac{\partial u}{\partial t})) - \frac{d\langle u \rangle}{dt} ((f'(u) \frac{\partial u}{\partial t}, 1)) \\ &\geq -c_0 \|\frac{\partial u}{\partial t}\|^2 - \frac{d\langle u \rangle}{dt} \frac{d}{dt} \int_{\Omega} F(u) dx \\ &= -c_0 \|\frac{\partial u}{\partial t}\|^2 - \frac{d}{dt} (\frac{d\langle u \rangle}{dt} \int_{\Omega} F(u) dx) + \frac{d^2 \langle u \rangle}{dt^2} \int_{\Omega} F(u) dx \end{aligned}$$

Set

$$\Lambda = \frac{1}{2} \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 - \frac{d\langle u \rangle}{dt} \int_{\Omega} F(u) dx$$

Thus :

$$\frac{d\Lambda}{dt} + \beta \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 + \left\| \frac{\partial \bar{u}}{\partial t} \right\|_V^2 \leq c_0 \left\| \frac{\partial u}{\partial t} \right\|^2 - \frac{d^2 \langle u \rangle}{dt^2} \int_{\Omega} F(u) dx$$

and

$$\frac{d\Lambda}{dt} + \beta \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 + \frac{1}{2} \left\| \frac{\partial \bar{u}}{\partial t} \right\|_V^2 \leq c \left(\left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 + \left| \frac{d\langle u \rangle}{dt} \right|^2 \right) - \frac{d^2 \langle u \rangle}{dt^2} \int_{\Omega} F(u) dx$$

Λ is bounded from below :

$$\Lambda \geq \frac{1}{2} \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 - c$$

The last two terms in the right-hand side are bounded from above :

$$\frac{d\Lambda}{dt} + \beta \left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 + \frac{1}{2} \left\| \frac{\partial \bar{u}}{\partial t} \right\|_V^2 \leq c \left(\left\| \frac{\partial \bar{u}}{\partial t} \right\|_{-1}^2 + \langle u_0 \rangle^2 + E(0)^2 + 1 \right)$$

→ Uniform Gronwall's lemma ($|\langle \frac{\partial u}{\partial t} \rangle| = \beta |\langle u \rangle| \leq \beta$)

Proposition : We assume that $2 \leq p < +\infty$, when $n = 2$, and $2 \leq p \leq 6$, when $n = 3$. Then, the solution u further satisfies

$$\|f(u)\|_{L^\infty(r,t;L^p(\Omega))} \leq c,$$

$$\|u(t)\|_{W^{2,p}(\Omega)} \leq c,$$

$\forall t \geq r, r > 0$ given, where the constant c depends on the $H^1(\Omega)$ -norm of u_0 .

Consider the elliptic equation

$$-\Delta u + f_1(u) = \mu + c_0 u, \quad \frac{\partial u}{\partial \nu} = 0 \text{ on } \Gamma$$

It suffices to prove that $\mu \in L^\infty(r, t; H^1(\Omega))$

Note that $\bar{\mu} \in L^\infty(r, t; H^1(\Omega))$

Multiply

$$-\Delta \bar{u} + \overline{f(u)} = \bar{\mu}$$

by $\overline{f(u)}$:

$$\|\overline{f(u)}\|^2 \leq c(\|\nabla u\|^2 + \|\nabla \mu\|^2)$$

$$\rightarrow \overline{f(u)} \in L^\infty(r, t; L^2(\Omega))$$

Finally : $\langle \mu \rangle = \langle f(u) \rangle$

Proposition : We assume that $n = 1$. Then there exists $\delta \in (0, 1)$ depending on the $H^1(\Omega)$ -norm of u_0 such that

$$\|u(t)\|_{L^\infty(\Omega)} \leq 1 - \delta, \quad t \geq r,$$

$r > 0$ given.

Proposition : We assume that $n = 2$. Then, the following holds for every $t \geq r$, $r > 0$ given, and for every $p \in \mathbb{N}$:

$$\|f'(u)\|_{L^p(\Omega \times (r, t))} \leq c,$$

where the constant c depends on p .

Proposition : We assume that $n = 2$. Then, the weak solution u further satisfies

$$\frac{\partial u}{\partial t} \in L^\infty(r, +\infty; H) \cap L^2(r, T; H^2(\Omega)),$$

$\forall r < T, r > 0$ and $T > 0$ given.

Theorem : We assume that $n = 2$. Then, there exists $\delta \in (0, 1)$ depending on the $H^1(\Omega)$ -norm of u_0 such that

$$\|u(t)\|_{L^\infty(\Omega)} \leq 1 - \delta, \quad t \geq r,$$

$r > 0$ given.

Remarks :

a) Strict separation property : existence of finite-dimensional attractors, convergence to steady states

b) To avoid differentiating with respect to time : quotient differences

Useful when one cannot differentiate with respect to time (e.g. : coupling with the incompressible Navier-Stokes equations)